

Non-self-adjoint scattering on graphs

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Scattering theory for self-adjoint operators

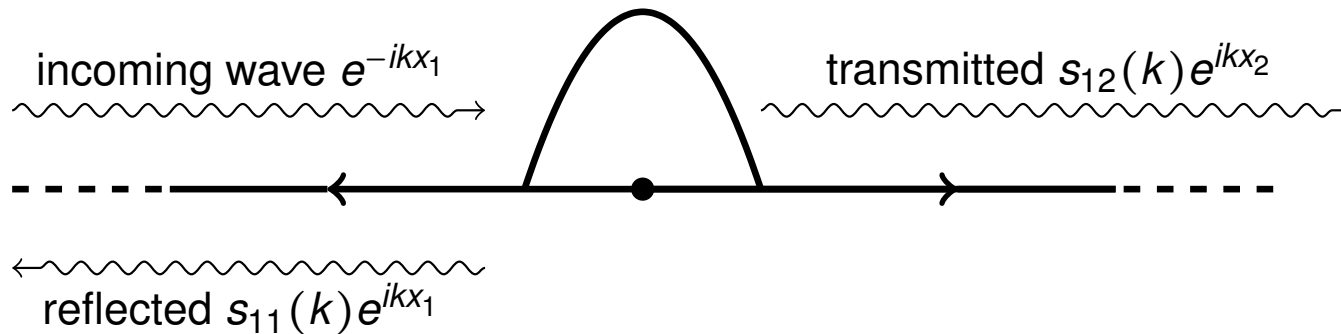
Schrödinger operators

$$\mathcal{H}_0 = -\frac{d^2}{dx^2} \quad \text{“free”}$$

$$\mathcal{H} = -\frac{d^2}{dx^2} + V \quad \text{“perturbed”}$$

Scattering matrix

$$S(k) = \begin{bmatrix} s_{11}(k) & s_{12}(k) \\ s_{21}(k) & s_{22}(k) \end{bmatrix}, \quad k > 0$$



- ▶ e.g. $V \in C_0^\infty(\mathbb{R}; \mathbb{R})$
- ▶ $\mathcal{H} = \mathcal{H}^*$ in $L^2(\mathbb{R})$
- ▶ $(e^{it\mathcal{H}})_{t \in \mathbb{R}}$ unitary group

- ▶ $S(k)$ for $k > 0$ unitary
- ▶ $S(\cdot)$ describes scattering states of \mathcal{H} by scattering states of $\mathcal{H}_0 = -\frac{d^2}{dx^2}$

Scattering theory for self-adjoint operators

Wave operators

$$W_+(\mathcal{H}, \mathcal{H}_0) := s - \lim_{t \rightarrow +\infty} e^{it\mathcal{H}} e^{-it\mathcal{H}_0} P_{\mathcal{H}_0}^{ac}$$

$$W_-(\mathcal{H}, \mathcal{H}_0) := s - \lim_{t \rightarrow -\infty} e^{it\mathcal{H}} e^{-it\mathcal{H}_0} P_{\mathcal{H}_0}^{ac}$$

Scattering operators

$$S := W_+^*(\mathcal{H}, \mathcal{H}_0) W_-(\mathcal{H}, \mathcal{H}_0)$$

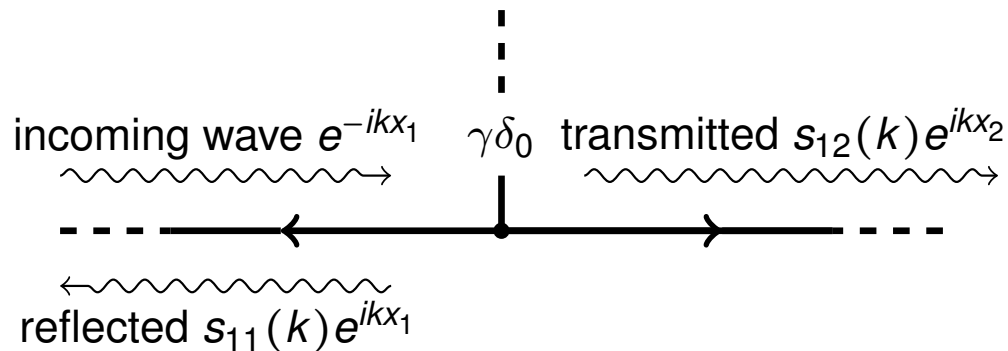
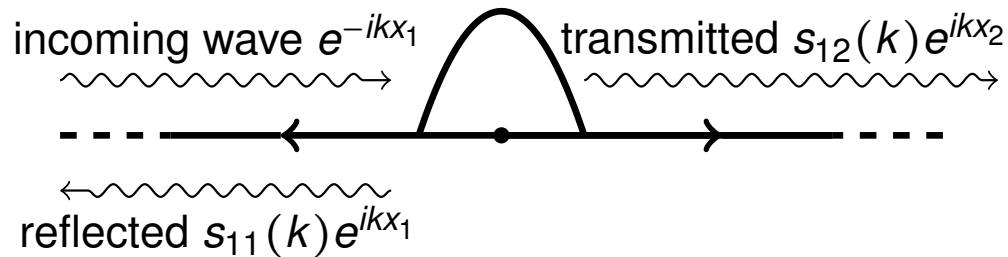
Absolutely continuous spectrum

- ▶ For $\mathcal{H}_0, \mathcal{H}$ self-adjoint absolutely continuous spectrum defined via spectral measure
- ▶ Absolutely continuous part defined via image of spectral projection

Existence and Completeness of W_{\pm}

- ▶ If W_{\pm} exist and surjective, then absolutely continuous part of \mathcal{H}_0 and \mathcal{H} unitarily equivalent
- ▶ Scattering theory \approx perturbation theory for absolutely continuous spectrum

Scattering theory for self-adjoint operators: point interaction



Point interactions = boundary conditions

$$D(\Delta(A, B)) := \{\psi \in H^2(0, \infty; \mathbb{C})^2 : A\psi(0) + B\partial_x\psi(0) = 0\},$$

$$\Delta(A, B)\psi := \partial_x^2\psi$$

Generalized eigenfunctions

Find $S(k)$ such that

$$\psi(x, k) = e^{-ikx} + S(k)e^{ikx}$$

satisfies

$$A\psi(0, k) + B\partial_x\psi(0, k) = 0$$

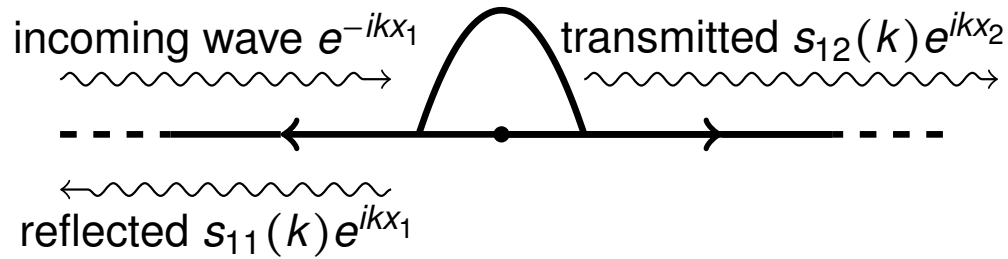
then

$$S(k) = -(A + ikB)^{-1}(A - ikB)$$

Scattering matrix

- ▶ $S(k)$ for $k > 0$ unitary if and only if $-\Delta(A, B)$ self-adjoint
- ▶ $S(\cdot)$ defines scattering operator $S = W_+^* W_-$
- ▶ $\psi(\cdot, k)$ span absolutely continuous part of $-\Delta(A, B)$

Scattering theory for dissipative operators



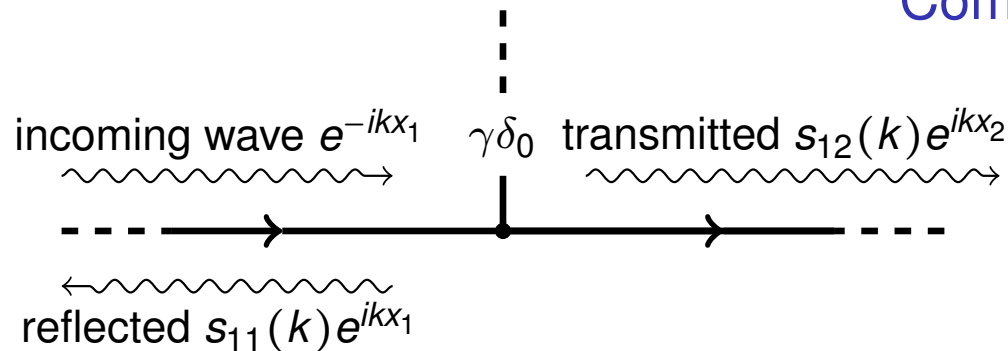
Optical model

$$\mathcal{H} = \mathcal{H}_0 + V + iC^*C,$$

$$V = V^* \text{ and } C \in \mathcal{L}(X)$$

- ▶ Davies (1978, ...)
- ▶ Faupin, Fröhlich et al. (2018, ...)

Complex δ -coupling

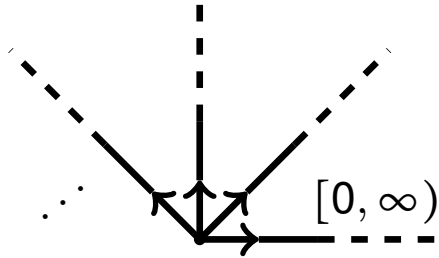


$$\mathcal{H} = \mathcal{H}_0 + \alpha\delta_0,$$

$$\alpha \in \{\operatorname{Re} \alpha \leq 0, \operatorname{Im} \alpha < 0\}$$

- ▶ Kadowaki et al. (2003, ...)
- ▶ Behrndt, Malamud, Neidhardt (2007, ...)

Scattering theory for non-self-adjoint operators on star graphs



Spectral singularities

$$\det(A + ikB) = 0$$

- ▶ $\text{Im } k > 0$, then $-k^2$ eigenvalue
- ▶ otherwise k spectral singularity

Generalized eigenfunctions

Find $S(k)$ such that

$$\psi(x, k) = e^{-ikx} + S(k)e^{ikx}$$

satisfies

$$A\psi(0, k) + B\partial_x\psi(0, k) = 0$$

then

$$S(k) = -(A + ikB)^{-1}(A - ikB)$$

Resolvent

Integral kernel for $\text{Im } k > 0$

$$r_{A,B}(x, y; -k^2) = \frac{i}{2k} \text{diag}\{e^{ik|x_j - y_j|}\}_{j \in \mathcal{E}} + \frac{i}{2k} \text{diag}\{e^{ikx_j}\}_{j \in \mathcal{E}} S(k, A, B) \text{diag}\{e^{iky_j}\}_{j \in \mathcal{E}}$$

Model case on the half line

Operators and forms

Consider in $L^2(0, \infty; \mathbb{C})$ the operator

$$\Delta_\gamma u = \partial_x^2 u \quad \text{with} \quad D(\Delta_\gamma) := \{u \in H^2(0, \infty; \mathbb{C}) : \gamma u(0) + u'(0) = 0\}, \quad \gamma \in \mathbb{C} \cup \{\infty\}$$

The operator $-\Delta_\gamma$ is associated with the quadratic form

$$\begin{aligned} \delta_\gamma[u] &:= \int_0^\infty |\partial_x u(x)|^2 dx - \gamma |u(0)|^2, \quad u \in H^1(0, \infty; \mathbb{C}) \quad \text{if } \gamma \in \mathbb{C}, \\ \delta_\infty[u] &:= \int_0^\infty |\partial_x u(x)|^2 dx, \quad u \in H_0^1(0, \infty; \mathbb{C}) \quad \text{if } \gamma = \infty. \end{aligned}$$

Adjoint in $L^2(0, \infty; \mathbb{C})$

$$(\Delta_\gamma)^* = \Delta_{\bar{\gamma}}, \quad \gamma \in \mathbb{C} \cup \{\infty\}.$$

Spectrum and spectral singularities

$$\begin{aligned} \sigma_r(-\Delta_\gamma) = \emptyset, \quad \sigma_c(-\Delta_\gamma) = [0, \infty) \quad \text{and} \quad \sigma_p(-\Delta_\gamma) = \sigma_{pp}(-\Delta_\gamma) &= \begin{cases} \emptyset & \text{if } \operatorname{Re} \gamma \leq 0, \\ -\gamma^2 & \text{if } \operatorname{Re} \gamma > 0, \end{cases} \\ \text{spectral singularity} &= \begin{cases} -\gamma^2 & \text{if } \gamma \in \mathbb{R}, \\ \emptyset & \text{else,} \end{cases} \end{aligned}$$

Model case on the half line

C_0 -contraction semigroups

- (a) The operator $+i\Delta_\gamma$ is m-dissipative if and only if $\operatorname{Re}(-i\gamma) \leq 0$, that is, $\operatorname{Im}(\gamma) \leq 0$;
- (b) The operator $-i\Delta_\gamma$ is m-dissipative if and only if $\operatorname{Re}(+i\gamma) \leq 0$, that is, $\operatorname{Im}(\gamma) \geq 0$;
- (c) The operator $i\Delta_\gamma$ is skew-self-adjoint if and only if $\operatorname{Im}(\gamma) = 0$;
- (d) All $\pm i\Delta_\gamma$ generate (not necessarily) bounded C_0 -groups.

Pre-wave operators

$$W(-\Delta_\gamma, -\Delta_0, t) := e^{it\Delta_\gamma} e^{-it\Delta_0} \quad \text{for } t \in \mathbb{R},$$

$$W(-\Delta_0, -\Delta_\gamma, t) := e^{it\Delta_0} e^{-it\Delta_\gamma} P_\rho^\perp \quad \text{for } t \in \mathbb{R},$$

Absolutely continuous subspace

If $-i\mathcal{H}$ dissipative¹, then

$$H_{ac}(\mathcal{H}) := \overline{M(\mathcal{H})}$$

where

$$M(\mathcal{H}) := \left\{ u \in H, \exists c_u > 0, \forall v \in H, \int_0^\infty | \langle e^{-it\mathcal{H}} u, v \rangle |^2 dt \leq c_u \|v\|^2 \right\}.$$

¹compare Davies (1978)

Model case on the half line: Existence of wave operators

Resolvent difference

$$r_{A,B}(x, y; -k^2) = \frac{i}{2k} e^{ik|x-y|} + \frac{i}{2k} e^{ikx} S(k, A, B) e^{iky}$$

Therefore,

$$r_{0,1}(x, y; -k^2) - r_{A,B}(x, y; -k^2) = \frac{i}{2k} e^{ikx} [1 - S(k, A, B)] e^{iky}$$

and

$$\frac{i}{2k} e^{ikx} e^{iky} = \frac{1}{2} (r_{0,1}(x, y; -k^2) - r_{1,0}(x, y; -k^2))$$

EnB method

Splitting into “Ingoing” and “outgoing” spaces using dilation operator

$$\frac{1}{2i} (x\partial_x + \partial_x x)$$

... \Rightarrow Existence wave operators

Cook-Kuroda method

Difference of resolvents trace class

... \Rightarrow Existence wave operators

Extension to self-adjoint system

Extension theory for symmetric operators with finite deficiency indices

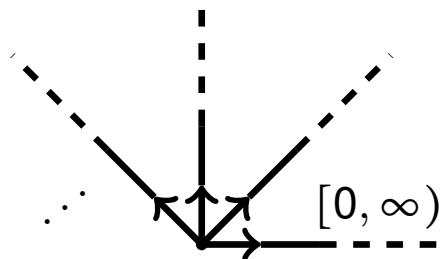
... \Rightarrow Existence wave operators

Absolutely continuous subspace

Replace semigroup by resolvents

... $\Rightarrow H_{ac}(i\Delta_\gamma) = H_p((i\Delta_\gamma)^*)^\perp$

Scattering theory for non-self-adjoint operators on star graphs



Assumption

For orthogonal projection P ,

$$A = L + P \quad \text{and} \quad B = P^\perp, \quad \text{where}$$

$$L = P^\perp L P^\perp \quad \text{and} \quad L \text{ diagonalizable}$$

$$\text{Im}(L) \leq 0$$

Generalized eigenfunctions

Find $S(k)$ such that

$$\psi(x, k) = e^{-ikx} + S(k)e^{ikx}$$

satisfies

$$A\psi(0, k) + B\partial_x\psi(0, k) = 0$$

then

$$S(k) = -(A + ikB)^{-1} (A - ikB)$$

Wave operator

Then one has the existence of

$$s - \lim_{t \rightarrow -\infty} W(-\Delta(0, 1), -\Delta(A, B), t)$$

$$= s - \lim_{t \rightarrow -\infty} e^{it\Delta(0,1)} e^{-it\Delta(A,B)} P_p^\perp$$

$$s - \lim_{t \rightarrow +\infty} W(-\Delta(A, B), -\Delta(0, 1), t)$$

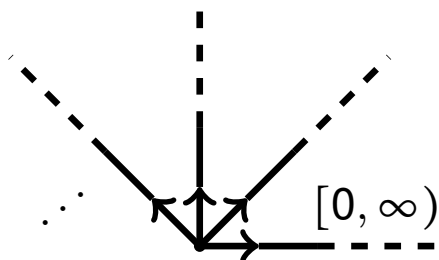
$$= s - \lim_{t \rightarrow +\infty} e^{it\Delta(A,B)} e^{-it\Delta(0,1)}$$

and on the half line

$$s - \lim_{t \rightarrow +\infty} W(-\Delta_\gamma, -\Delta_0, t)f = \int_0^\infty \psi(\cdot, k) \hat{f}(k) \frac{dk}{2},$$

$$\hat{f}(k) = \frac{2}{\pi} \int_0^\infty \cos(xk) f(x) dx$$

Scattering theory for non-self-adjoint operators on star graphs



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Spectral singularities

$$\det(A + ikB) = 0$$

- ▶ $\text{Im } k > 0$, then $-k^2$ eigenvalue
- ▶ otherwise k spectral singularity

Conjecture

Wave operators

$$s - \lim_{t \rightarrow -\infty} W(-\Delta(0, 1), -\Delta(A, B), t)$$

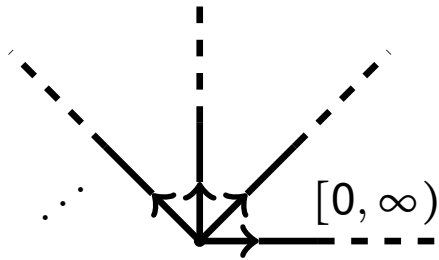
$$= s - \lim_{t \rightarrow -\infty} e^{it\Delta(0,1)} e^{-it\Delta(A,B)} P_p^\perp$$

$$s - \lim_{t \rightarrow +\infty} W(-\Delta(A, B), -\Delta(0, 1), t)$$

$$= s - \lim_{t \rightarrow +\infty} e^{it\Delta(A,B)} e^{-it\Delta(0,1)}$$

(asymptotically) complete if and only if there is **no** spectral singularity

Scattering theory for non-self-adjoint operators on star graphs



Assumption

For orthogonal projection P ,

$$A = L + P \quad \text{and} \quad B = P^\perp, \quad \text{where}$$

$$L = P^\perp L P^\perp \quad \text{and} \quad L \text{ diagonalizable}$$

$$\text{Im}(L) \leq 0$$

Spectral singularities

$$\det(A + ikB) = 0$$

- ▶ $\text{Im } k > 0$, then $-k^2$ eigenvalue
- ▶ otherwise k spectral singularity

Thanks for your attention!

Conjecture

Wave operators

$$s - \lim_{t \rightarrow -\infty} W(-\Delta(0, 1), -\Delta(A, B), t)$$

$$= s - \lim_{t \rightarrow -\infty} e^{it\Delta(0,1)} e^{-it\Delta(A,B)} P_p^\perp$$

$$s - \lim_{t \rightarrow +\infty} W(-\Delta(A, B), -\Delta(0, 1), t)$$

$$= s - \lim_{t \rightarrow +\infty} e^{it\Delta(A,B)} e^{-it\Delta(0,1)}$$

(asymptotically) complete if and only if there is **no** spectral singularity